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## Fast Transient Fluorescence (FTRF) Technique for Swelling of Gels in Solvent Mixtures

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### ABSTRACT

Strobe Master System (SMS), is introduced for studying THE swelling of disc shaped gels in chloroform–heptane mixtures. Gels were prepared by free radical copolymerization of methyl (methacrylate) (MMA) and ethylene glycol dimethacrylate (EGDM). Pyrene (P) was introduced as a fluorescence probe during polymerization and the existence of P were measured using (SMS) during *in situ* swelling processes. Various percents of chloroform in chloroform–heptane mixtures were used as a swelling agent. For the duration,  $\tau_2$  of P in the gel were found to be decreased as swelling proceeded. Li–Tanaka's equation was employed to determine the time constants,  $\tau_c$  and cooperative diffusion coefficients,  $D_c$  which were found to be decreased and increased respectively by increasing the chloroform content in the solvent mixtures.

### INTRODUCTION

Fluorescence dyes can be used to study local environments, basically with two types of experiments. When the dye is simply added to the system as a small molecule, the dye is referred to as a probe which is available commercially. Consequently, such experiments are easy to carry out, but are often difficult to interpret because one has to know where the

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dye is located in the system. If one can prepare an experiment which allows the dye to be attached covalently to a specific component of a system such as a polymer chain segment, the dyes are referred to as labels. The question can be raised whether the presence of the dye disturbs the system or disturbs its own local environments in the system. Disturbances are most common where high dye concentration leads to aggregation, and in crystalline systems where the order in the system can be affected by the dye. Disturbances are much less likely when the fluorescent dye is incorporated into an amorphous fluid or glassy phase.

Volume phase transitions in gels may occur from dry to swollen states either continuously, or by sudden jumps from one state to the other.<sup>[1,2]</sup> The equilibrium swelling of gels in solvent has been studied extensively.<sup>[3–5]</sup> The swelling process of chemically cross-linked gels can be understood by considering the osmotic pressure vs. the restraining elastic force.<sup>[6–8]</sup> The total free energy of a chemical gel consists of bulk and shear energies. In fact, in a swollen gel, the bulk energy can be characterized by the osmotic bulk modulus  $K$ , which is defined in terms of the swelling pressure, and the volume fraction of polymer at a given temperature. On the other hand, the shear energy that keeps the gel in shape can be characterized by the shear modulus  $G$ . The shear energy minimizes the nonisotropic deformations in the gel. The theory of swelling kinetics for a spherical chemical gel was first developed by Tanaka and Filmore,<sup>[9]</sup> where the assumption is made that the shear modulus  $G$  is negligible compared to the osmotic bulk modulus. Peters and Candau<sup>[10]</sup> derived a model for the kinetics of swelling in spherical and cylindrical gels by assuming non-negligible shear-modulus and Li and Tanaka<sup>[6]</sup> developed a model where the shear modulus plays an important role that keeps the gel in shape due to coupling of any change in different directions. This model predicts that the geometry of the gel is an important factor, and swelling is not a pure diffusion process.

Many different experimental techniques have been used to study the kinetics of swelling and shrinking of chemical and physical gels among which are neutron scattering,<sup>[11]</sup> quasielastic light-scattering<sup>[10]</sup> macroscopic experiments<sup>[7]</sup> and *in situ* interferometric<sup>[12]</sup> measurements. Time resolved and steady-state fluorescence techniques were employed to study isotactic polystyrene in its gel state<sup>[13]</sup> where excimer spectra were used to monitor the behavior of two different conformations in the gel state of polystyrene. A pyrene derivative was used as a fluorescence molecule for monitoring the polymerization, aging and drying of aluminosilicate gels.<sup>[14]</sup> These results were interpreted in terms of chemical changes occurring during the sol–gel transition and the interactions between the chromophores and the sol–gel matrix.

*In situ* observations of the sol–gel transition in free-radical crosslinking copolymerization, using the steady-state fluorescence technique, were reported.<sup>[15,16]</sup> The same technique was also applied for studying swelling and drying kinetics in disc shaped gels.<sup>[17,18]</sup> Recently, fast transient fluorescence technique (FTRF) was used to monitor the swelling of poly(methyl methacrylate) (PMMA) gels.<sup>[19,20]</sup>

In this work, swelling of gels formed by free radical copolymerization (FCC) of methyl-methacrylate (MMA) and ethylene glycol dimethacrylate (EGDM) was studied in chloroform–heptane mixtures using the FTRF technique. Lifetimes of pyrene (P) which is embedded in the gel were followed during the *in situ* swelling processes. A Strobe Master System (SMS) was used for duration measurements of P in the gel. Duration measurements with SMS takes a much shorter time than is required in single photon counting systems and phase instruments. This advantage of SMS allows one to make

hundreds of measurements during the swelling process of gels. For this reason, we named this technique Fast Transient Fluorescence (FTRF), which gives us many advantages compared to other existing measuring techniques. It is observed that as the gel swells the existence of P in the gel decreases which can be modeled by the using low quenching Stern–Volmer equation. Cooperative diffusion coefficients,  $D_c$  were determined in respect to the chloroform content in the chloroform–heptane mixtures by employing the Li–Tanaka equation.

### KINETICS OF SWELLING

It has been suggested<sup>[6]</sup> that the kinetics of swelling and shrinking of a polymer network or gel should obey the following relation:

$$\frac{W(t)}{W_\infty} = 1 - \sum_{n=1}^{\infty} B_n e^{-t/\tau_n} \quad (1)$$

Where  $W(t)$  and  $W_\infty$  are the swelling or solvent uptake at time  $t$  and infinite equilibrium, respectively.  $W(t)$  can also be considered as volume difference of the gel at the times  $t$  and zero. Each component of the displacement vector of a point in the network from its final equilibrium location, after the gel is fully swollen, decays exponentially with a time constant  $\tau_n$  which is independent of time  $t$ . Here  $B_n$  is given by following relationship<sup>[6]</sup>:

$$B_n = \frac{2(3 - 4R)}{\alpha_n^2 - (4R - 1)(3 - 4R)} \quad (2)$$

where  $R$  is defined as the ratio of the shear and the longitudinal osmotic modulus,  $R = G/M$ . The longitudinal osmotic modulus,  $M$  is a combination of shear,  $G$  and osmotic bulk modulus,  $K$ ,  $M = K + 4G/3$ , and  $\alpha_n$  are the roots of the Bessel function given as a function of  $R$  as follows:

$$R = \frac{1}{4} \left[ 1 + \frac{\alpha_n J_0(\alpha_n)}{J_1(\alpha_n)} \right] \quad (3)$$

Here  $J_0$  and  $J_1$  are the Bessel functions of the zeroth and first order.

In Eq. (1),  $\tau_n$  is inversely proportional to the collective cooperative diffusion coefficient  $D_c$  of a gel disc at the surface and is given by the relationship<sup>[7]</sup>

$$\tau_n = \frac{3a^2}{D_c \alpha_n^2} \quad (4)$$

Here, the diffusion coefficient  $D_c$  is given by  $D_c = M/f = (K + 4G/3)/f$ ,  $f$  is the friction coefficient describing the viscous interaction between the polymer and the solvent, and  $a$  represent half of the disc thickness in the final infinite equilibrium which can experimentally be determined.

The series given by Eq. (1) is convergent. The first term of the series expansion is dominant at large  $t$ , which correspond to the final stage of the swelling. If  $n > 1$ ,  $\alpha_n$  increases and  $\tau_n$  decreases very rapidly [see Eq. (4)]. Therefore, the kinetics of swelling in

the limit of large  $t$ , or if  $\tau_1$  is much larger than the rest of  $\tau_n$ , all higher order terms ( $n \geq 2$ ) in Eq. (1) can be neglected, and the swelling and shrinking can be represented by first order kinetics.<sup>[6]</sup> In this case, Eq. (1) simplifies to:

$$\frac{W(t)}{W_\infty} = 1 - B_1 e^{-t_s/\tau_c} \quad (5)$$

Equation 5 allows us to determine the parameters  $B_1$  and  $\tau_c$ .

It is important to note that Eq. (5) satisfies the following equation:

$$\frac{dW(t)}{dt} = \frac{1}{\tau_c} (W_\infty - W) \quad (6)$$

which proves the mentioned the first order kinetics. The higher order terms ( $n \geq 2$ ) can be considered as fast decaying perturbation to the first order kinetics of swelling in the limit of large  $t$ .

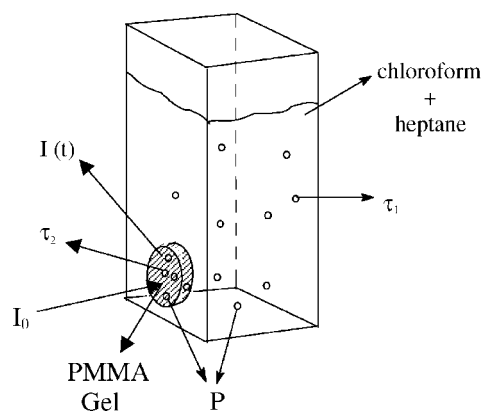
## EXPERIMENTAL

EGDM has been commonly used as a crosslinker in the synthesis of polymeric networks. The monomers MMA (Merck) and EGDM (Merck) were freed from the inhibitor by shaking with a 10% aqueous KOH solution, washing with water, and drying over sodium sulfate. They were then distilled under reduced pressure over copper chloride. The swelling agent chloroform (Merck), was distilled twice over sodium.

The radical copolymerization of MMA and EGDM was performed at 65°C in the presence of 2,2'-azobisisobutyronitrile (AIBN) (0.26 wt%) as an initiator. P was added as a fluorescence probe during the gelation process. The sample was deoxygenated by bubbling nitrogen through the solution for 10 min, then the radical copolymerization of MMA and EGDM was performed at  $65 \pm 2^\circ\text{C}$ . The EGDM content was kept at 0.035 vol%, and the P concentration was taken as  $4 \times 10^{-4}$  M. After completed gelation, the cylindrical gel sample was dried under vacuum and cut into the disc shaped gels for the swelling experiments.

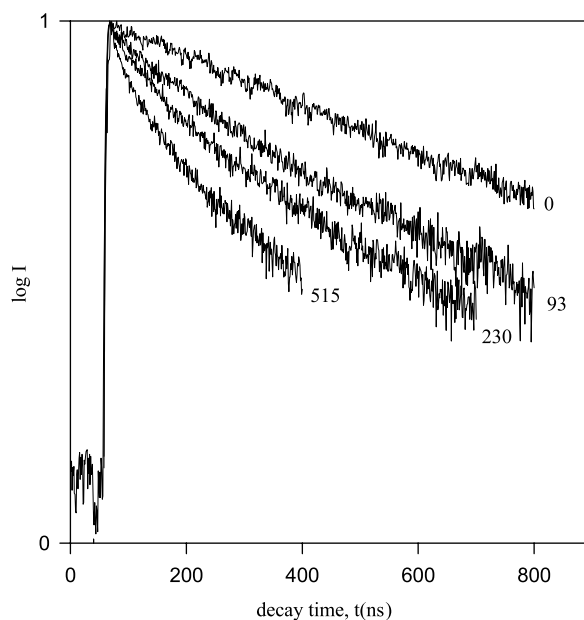
Fluorescence decay experiments were performed using a Photon Technology International (PTI) Strobe Master System (SMS). In the strobe, or pulse sampling technique<sup>[21,22]</sup> the sample is energized by a pulsed light source. The name arises from the Photo Multiplier Tube (PMT) that is gated or *strobed* by a voltage pulse that is synchronized with the pulsed light source. The intensity of the fluorescence emission is measured in a very narrow time window on each pulse and saved in a computer. The time window is moved to a larger time after each pulse. The strobe has the effect of turning of the PMT and measuring the emission intensity over a very short time window. When the data have been sampled over an appropriate range of time, a decay curve of fluorescence intensity as a function of time can be established.

The *in situ* swelling experiment was carried out in the SMS of PTI, by employing a pulsed lamp source (0.5 atm of  $\text{N}_2$ ). Pyrenes in the gel sample were excited at 345 nm and fluorescence decay curves were obtained at 395 nm during the *in situ* swelling experiment which was performed at room temperature. A disc shaped gel sample was placed in

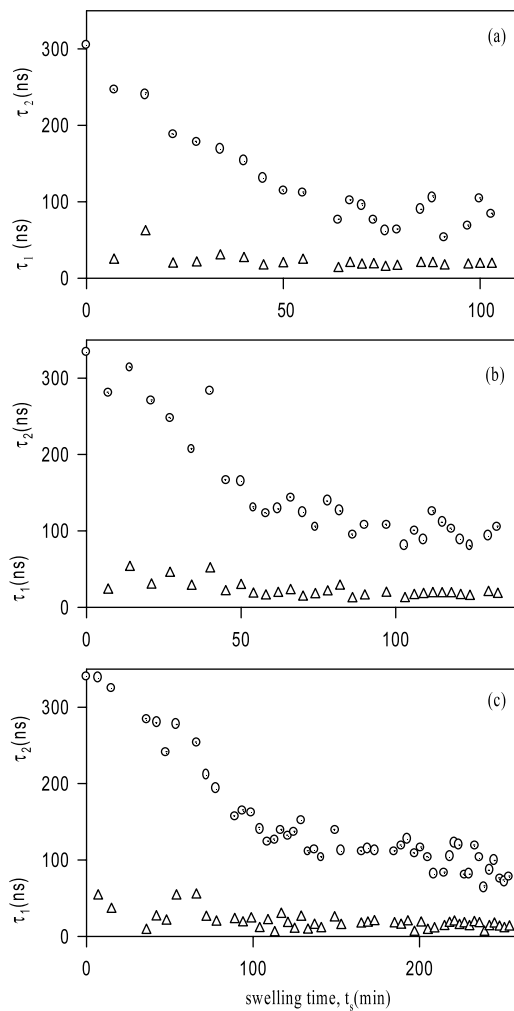


**Figure 1.** Cartoon representation of the gel position in the fluorescence cell.  $I_0$  and  $I$  are the incident and fluorescence light beams respectively. Small dots present the Pyrene molecules.  $\tau_2$  and  $\tau_1$  are the pyrene lifetimes inside and outside the gel.

a 1 cm  $\times$  1 cm quartz cell filled with chloroform–heptane mixtures for swelling process. The position of the disc shaped gel and the incident and fluorescence light beams are shown in Fig. 1 for illustration. The fluorescence decay data were collected over three decades in time and fitted by nonlinear least squares of exponential functions using a deconvolution



**Figure 2.** Fluorescence profiles at various swelling steps. The number on each curve present the swelling time in minutes.

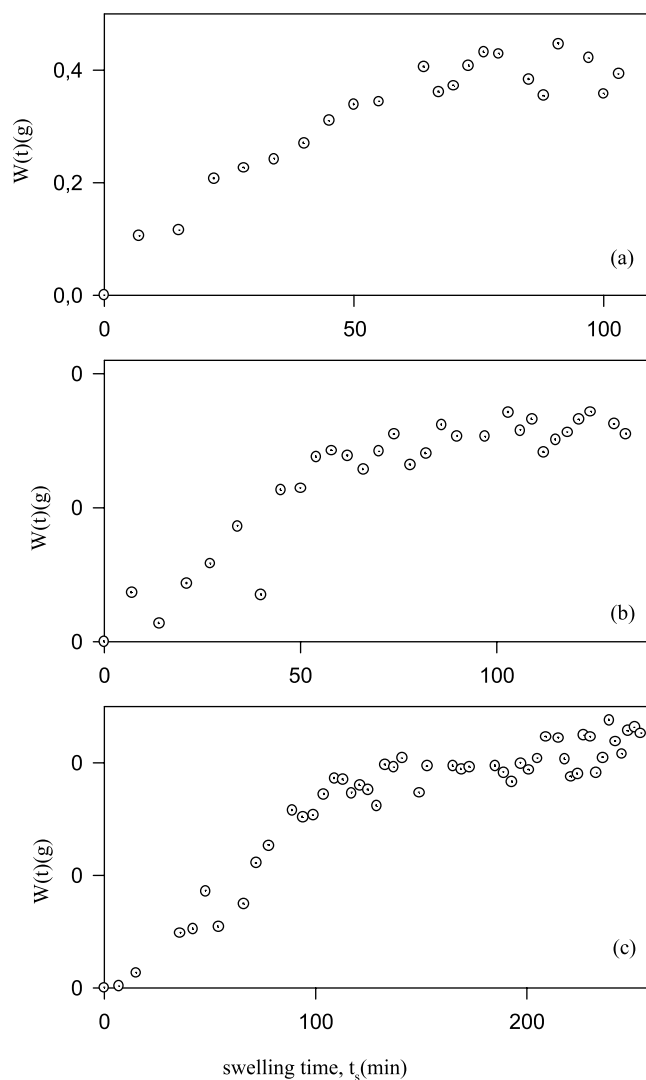


**Figure 3.** The plot of the measured  $\tau_1$  and  $\tau_2$  values versus swelling time,  $t_s$  for a—100, b—90, c—80% of chloroform contents.

method with a dry gel as a scatterer standard. The quality of the data fit to the model is checked by  $\chi^2$  ( $\chi^2 \leq 1.10$ ), the distribution of the weighted residuals and the autocorrelation of the residuals.

## RESULTS AND DISCUSSIONS

Figure 2 presents the fluorescence decay profiles at various swelling stages (0, 93, 230 and 515 min). As the swelling time,  $t_s$  is increased, the excited pyrenes decay faster and



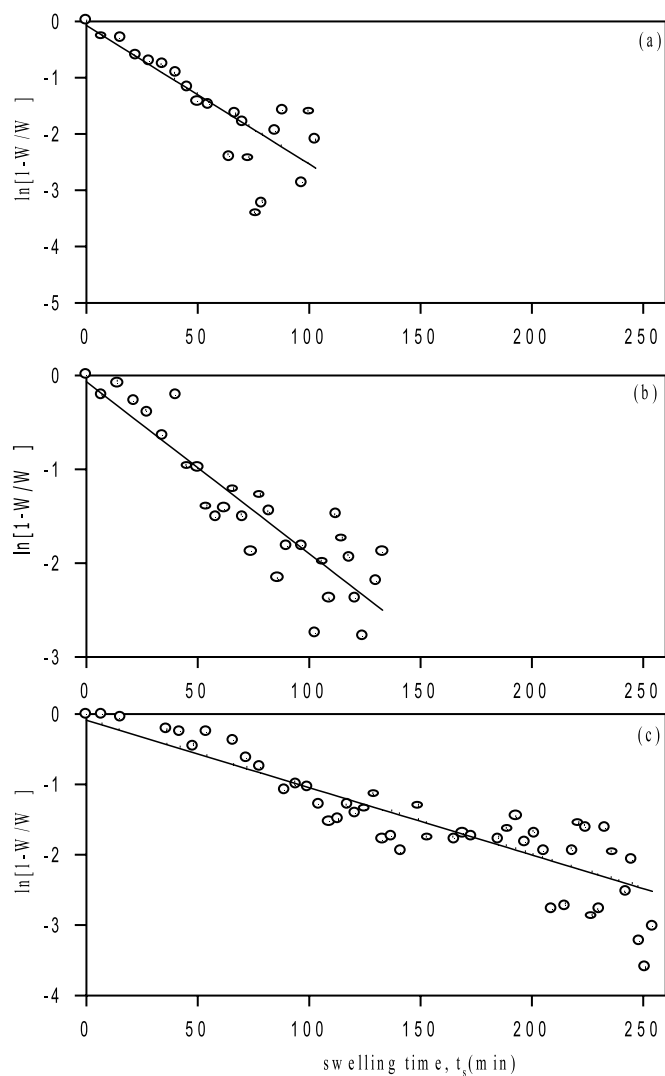
**Figure 4.** Plot of the solvent uptake,  $W$  vs. swelling time,  $t_s$  for a—100, b—90, c—80% of chloroform contents.

faster. It indicates that as the solvent uptake is increased, the quenching of excited pyrenes also increase. In order to probe the swelling process during solvent uptake, the fluorescence decay curves were measured and fitted to the sum of two exponentials:

$$I(t) = A_1 e^{-t/\tau_1} + A_2 e^{-t/\tau_2} \quad (7)$$

where  $\tau_1$  and  $\tau_2$  are the pyrene lifetimes outside and inside of the gel.  $A_1$  and  $A_2$  are the corresponding amplitudes of the decay curves. Here, the role of the solvent in quenching is to add the quasi-continuum of states needed to satisfy energy resonance





**Figure 5.** Fit of the data in Fig. 4 to Eq. (12). The slope of the curve produced  $\tau_c$  values which are listed in Table 1.

conditions, i.e., the solvent acts as an energy sink for rapid vibrational relaxation which occurs after the rate limiting transition from the initial state. Birks *et al.* studied the influence of solvent viscosity on the fluorescence characteristics of pyrene solutions in various solvents and observed that the monomer internal quenching rate is affected by the solvent quality.<sup>[23]</sup>

The measured  $\tau_1$  and  $\tau_2$  values for the gel samples swollen in various chloroform–heptane mixtures are plotted versus swelling time,  $t_s$  in Fig. 3. The  $\tau_1$  values did not change much, but the  $\tau_2$  values decreased as  $t_s$  is increased. The decrease in  $\tau_2$  arises from solvent

**Table 1.**

% Chloroform	$\Delta m$ (g)	$a_\infty$ (cm)	$\tau_c$ (s)	$B_1$	$\alpha_1$	$D_c$ ( $\times 10^{-5} \text{ cm}^2 \text{ s}^{-1}$ )
100	0,40	0.29	2460	0.94	0.87	3.38
95	0,36	0.28	5040	0.97	0.64	2.95
90	0,16	0.27	3276	0.94	0.87	2.20
85	0,14	0.31	4200	0.93	0.95	1.96
80	0,21	0.3	6240	0.91	1	1.08
70	0,16	0.34	8640	0.93	0.95	1.14
60	0,17	0.31	9120	0.85	1.27	0.49
40	0,13	0.29	12000	0.89	1.15	0.39
20	0,02	0.23	23370	0.9	1.14	0.12

Experimentally measured; maximum solvent uptake,  $\Delta m$ ; final disc thickness,  $a_\infty$ ; time constant,  $\tau_c$  and cooperative diffusion coefficient,  $D_c$ .

quenching while the gel swells. As the chloroform content is increased in solvent mixture,  $\tau_2$  decrease increasingly faster, i.e., in Fig. 3a,  $\tau_2$  decreases much more rapidly than in Fig. 3b,c. The effect of solvent quality can be interpreted by applying quenching mechanisms to the produced data where the Stern–Volmer type of quenching mechanism may be assumed for the fluorescence decay of P in the gel sample during the swelling process.<sup>[24]</sup> The following relationship for the existence in the gel was satisfied.

$$\tau_2^{-1} = \tau_{20}^{-1} + k_q[W] \quad (8)$$

where  $\tau_{20}$  is the duration of P in dry gel in which no quenching has taken place,  $k_q$  is the quenching rate constant and  $[W]$ , the solvent concentration in the gel for a certain solvent uptake. For low quenching efficiency, i.e.,  $\tau_{20}k_q[W] < 1$ , Eq. (8) becomes

$$\tau_2 \approx \tau_{20}(1 - \tau_{20}k_q[W]) \quad (9)$$

The solvent uptake,  $W$  is obtained from volume integration:

$$W = \int_{a_0}^{a_\infty} [W]dv \quad (10)$$

Where  $dv$  is the differential volume in the gel. The integration is taken from initial,  $a_0$  to final  $a_\infty$  thickness of the disc shaped gel. Performing the integration, the following relation is obtained:

$$W = \left(1 - \frac{\tau_2}{\tau_{20}}\right) \frac{\nu}{k_q \tau_{20}} \quad (11)$$

Here,  $\nu$  is the swollen volume of the gel, which can be measured experimentally.  $k_q$  was obtained from separate measurements.

The plots of solvent uptake,  $W$ , for the gels swollen in 100, 90 and 80% chloroform contents are shown in Fig. 4a–c, respectively. These are typical solvent uptake curves following the Li–Tanaka Equation [Eq. (5)].

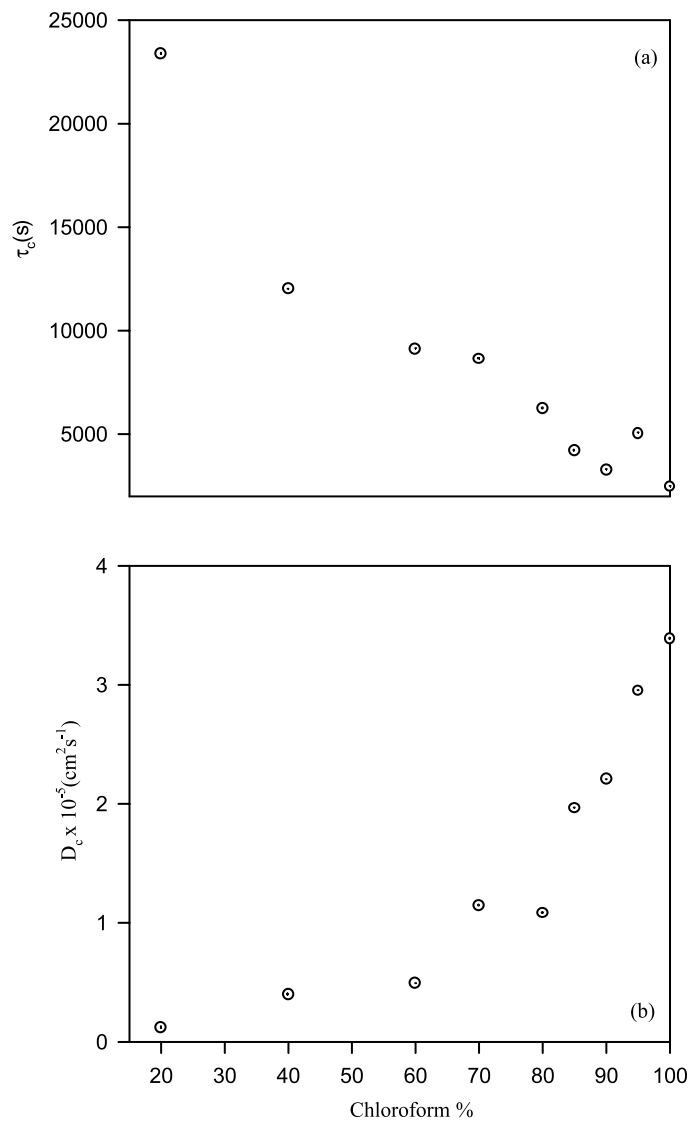
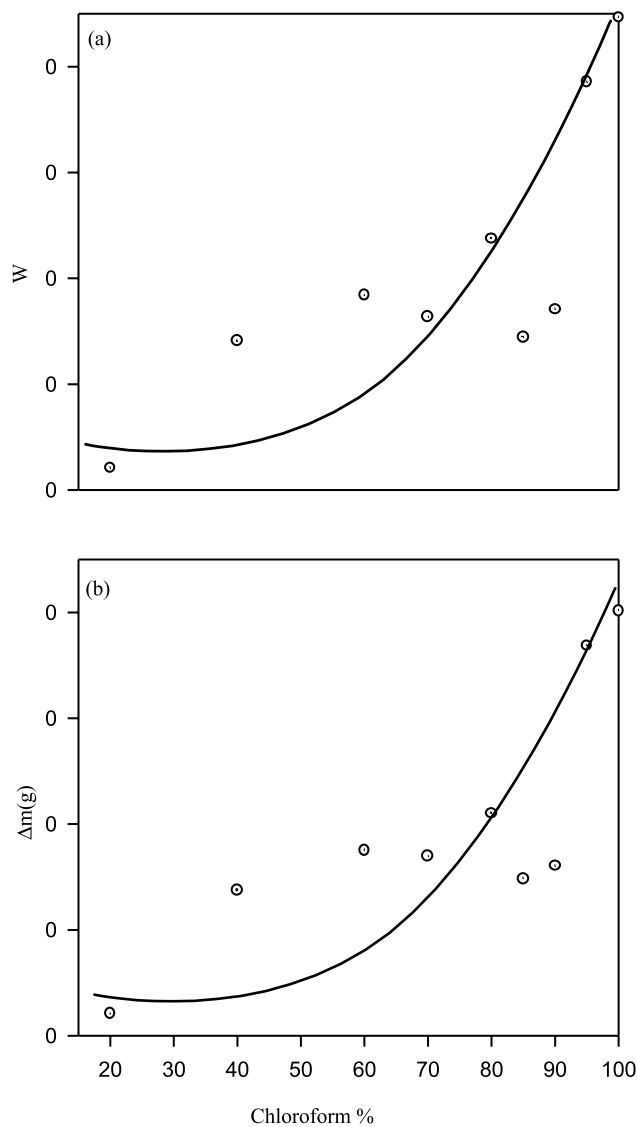


Figure 6. The plot of a— $\tau_c$  and b— $D_c$  values versus chloroform %.

The logarithmic data in Fig. 4 were fitted to the corresponding relationship of Eq. (5):

$$\ln \left( 1 - \frac{W}{W_\infty} \right) = \ln B_1 - \frac{t_s}{\tau_c} \quad (12)$$

The fits are presented in Fig. 5, from which  $B_1$  and  $\tau_c$  values were obtained. These are listed in Table 1. The plot of  $\tau_c$  vs. the chloroform percentage is shown in Fig. 6a. As



**Figure 7.** The plot of maximum solvent uptakes measured a—from existing measurements, b—from weighing measurements using micro-balance.

expected, the penetration of solvent molecules slows down as the heptane content is increased. A decrease in the solvent quality prevents the gel from swelling. Knowing  $B_1$ , one can obtain  $\alpha_1$  and from Eq. (4),  $D_c$ . These data are plotted in Fig. 6b vs. chloroform content. The behavior of  $D_c$  in Fig. 6b is evidence for gel segments which move faster in a thermodynamically high quality solvent, i.e., in chloroform containing solvents, gel segments diffuse considerably faster.

$W_{\infty}$  values could be obtained from the equilibrium regions in Fig. 4 and are plotted in Fig. 7a. Figure 7b presents the observed  $\Delta m$  values which were obtained by subtracting the final weight of the gel from its initial weight. In fact, the  $\Delta m$  and  $W_{\infty}$  values are the parameters similar to each other but measured by different techniques. As shown in Figs. 7a,b they behave quite similarly with regard of the chloroform content. Comparing Fig. 6b with Fig. 7, it is seen that in poor solvent quality (high heptane content) the gel does not swell much. The segments move rather slowly compared to those in gels swollen in the good solvent quality (high chloroform content).

### CONCLUSION

We showed that the FTRF technique can be used to measure cooperative diffusion coefficients,  $D_c$ , of the segment as a function of solvent quality during swelling a polymeric gel. One may speculate that measuring duration by FTRF in swelling gels may provide data which could be used for determining the rate of solvent penetration. In this paper, we introduced the FTRF method to study gel swelling in solvents of various quality. Solvents of poor solution quality penetrates the gel much more slowly than a solvent of good solution quality. Gels swell much faster and to a larger extent in high quality solvent than in bad quality solvents.

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